Quarkonium Production References

Geoff Bodwin

Pierre Artoisenet

Status of Theory Calculations & LHC Predictions
Workshop on Quarkonium Production
CERN, February 2010

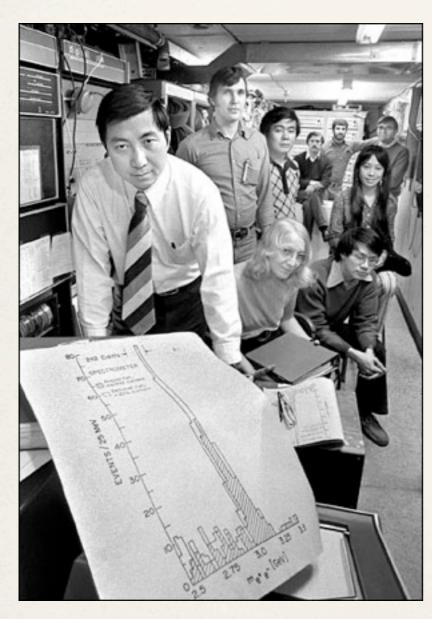
Stumbling towards a Theory of Quarkonium Production

Eric Braaten
The Ohio State University

Fermilab, May 21, 2010

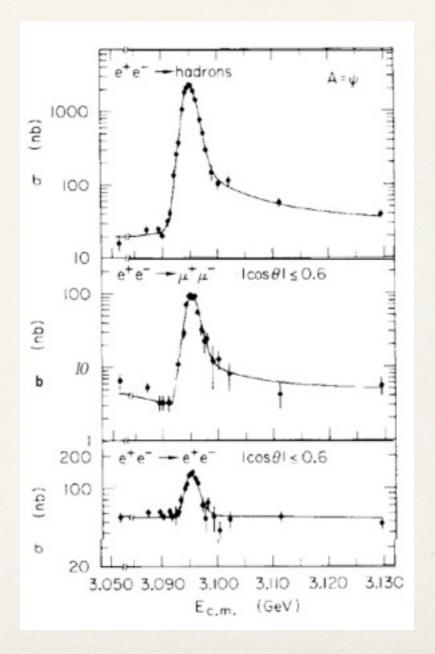
Discovery of Charmonium November 1974

p on Be target at Brookhaven



 $p N \rightarrow J/\psi + X, J/\psi \rightarrow e^+e^-$

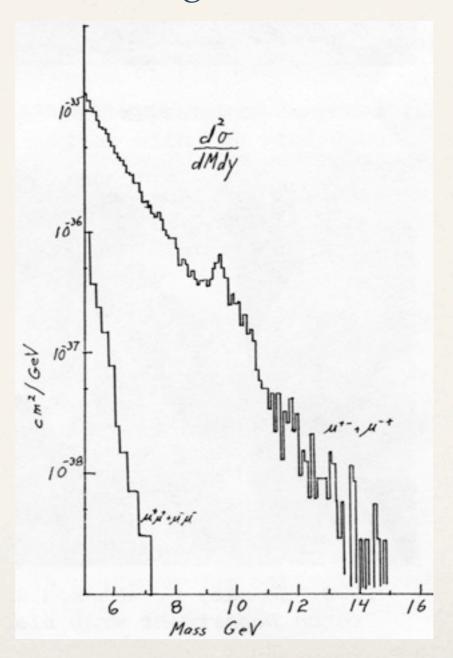
e+e⁻ annihilation at SLAC



 $e^+e^- \rightarrow J/\psi$, $J/\psi \rightarrow$ hadrons

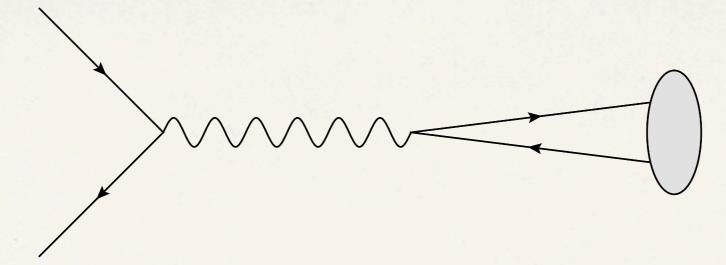
Discovery of Bottomonium August 1977

p on Be target at Fermilab



 $p N \rightarrow Y + X, Y \rightarrow \mu^+ \mu^-$

e^+e^- annihilation into J/ψ



- c c created by virtual photon
- rate determined by one constant f_{ψ}

$$\langle J/\psi|\bar{c}\gamma^{\mu}c|0\rangle = f_{\psi} \ \epsilon_{\psi}^{\mu}$$

- estimate using potential models: $f_{\psi} \propto R(0)$
- can be calculated using lattice QCD NRQCD, HPQCD, ...

Production of Charmonium using Hadrons

- How is the c \(\bar{c}\) pair created?
 What are the relevant parton processes?
 Can they be calculated using perturbative QCD?
- How does the $c\bar{c}$ pair bind to form charmonium? Can effects of binding be reduced to a few constants? Can they be calculated using lattice QCD?
- Possible answers: Color-singlet model (1976?)
 Color evaporation model (1977)
 NRQCD factorization (1995)

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Color-singlet Model

Ellis, Einhorn, Quigg 1976; Carlson and Suaya 1976; Kuhn 1980; Degrand, Toussaint 1980; Kuhn, Nussinov, Ruckl 1980; Wise 1980; Chang 1980; Baier, Ruckl 1981; Berger, Jones 1981

- *c* \bar{c} is created by parton collisions with negligible relative momentum
- $c \ \bar{c}$ can bind into charmonium only if it is created in same color/angular momentum as in charmonium

$$\frac{1}{3}$$
 S₁ for J/ ψ $\frac{1}{3}$ P_J for χ_{cJ}

• probability that $c \bar{c}$ binds into charmonium is determined by wavefunction near origin

$$\propto R(0)$$
 for J/ψ , η_c
 $\propto R'(0)$ for χ_{cl} , h_c

one constant for each multiplet can be determined from annihilation decays:

$$J/\psi \rightarrow e^+ e^-$$

$$\chi_{c0} \rightarrow \gamma \gamma$$

Color Evaporation Model

Fritzsch 1977; Halzen 1977

- $c \bar{c}$ pair is created by parton collisions with invariant mass below $D \bar{D}$ threshold (between $2m_c$ and $2m_D$)
- *c* \overline{c} c \overline{c} pair can bind into charmonium regardless of its color/angular momentum state
- probability that $c \bar{c}$ binds to form charmonium H is universal constant f_H for each multiplet

Color-singlet Model vs Color Evaporation Model

Applicability

CSM: exclusive and inclusive production definite predictions for polarization

CEM: only sufficiently inclusive production no polarization

Predictive power

CSM: one constant for each multiplet determined by annihilation decays

CEM: one constant for each multiplet adjustable parameters

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Color-singlet Model vs Color Evaporation Model

consistency

CSM: infrared divergences for P-waves

CEM: no infrared divergences

 $\chi_{cJ} \rightarrow q q g$ $b \rightarrow \chi_{cJ} + s + g$

perturbative corrections

CSM: separate NLO calculation for each process

CEM: can use NLO calculation for inclusive QQ

Nason, Dawson, Ellis 1988

• Dominant theoretical prejudice in early 1990's

CSM: can probably be extended

to a theory based on QCD

CEM: purely phenomenological model

Color-singlet Model vs Color Evaporation Model

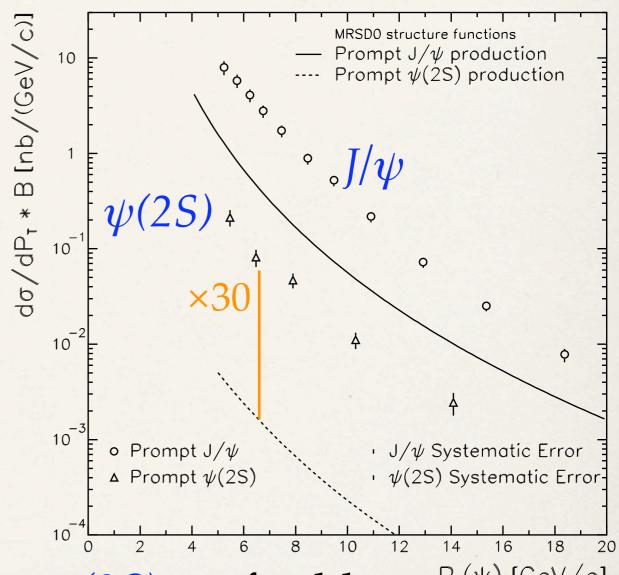
Experimental status in early 1990's

- Fixed target experiments $(pN, \pi N, \gamma N)$ feeddown to J/ψ from $\psi(2S)$, χ_{cJ} decays contributions from small $p_T \Rightarrow$ nonperturbative? large experimental errors roughly compatible with CSM or CEM
- $p \, \overline{p}$ collisions at the Tevatron feeddown to J/ψ from $\psi(2S)$, χ_{cJ} decays feeddown from B decays $p_T > 5 \, \text{GeV} \Rightarrow \text{perturbative?}$ production rates much larger than predicted by CSM?

Demise of Color-singlet Model

CDF collaboration 1997

• use vertex detector to remove *B* feeddown



- prompt J/ψ : complicated by $\psi(2S)$, χ_{cJ} feeddown [GeV/c]
- prompt $\psi(2S)$: 30 times larger than CSM prediction (in retrospect, compatible with CEM)

Nonrelativistic QCD

Caswell and Lepage 1986

- effective field theory for QQ sector of QCD at energies $<< m_Q$ from $Q\overline{Q}$ threshold
- in quarkonium, small velocity v is generated dynamically by balance between potential energy and kinetic energy

charmonium: $v^2 \approx 1/3$

bottomonium: $v^2 \approx 1/10$

nonperturbative effects
 can be organized according to their scaling with v

Nonrelativistic QCD (cont.)

Multipole expansion

- E1 transitions: $\Delta L = 1$, $\Delta S = 0$ amplitude ~ v
- M1 transitions: $\Delta L = 0$, $\Delta S = 1$ amplitude ~ v^2

Fock state expansion for quarkonium can be organized in powers of *v*

$$|J/\psi\rangle = \mathcal{O}(1) |c\bar{c}(\underline{1} {}^{3}S_{1})\rangle + \mathcal{O}(v) |c\bar{c}(\underline{8} {}^{3}P_{J}) + g\rangle + \mathcal{O}(v^{2})$$

Lattice NRQCD

Lepage et al. 1992

calculate properties of quarkonium nonperturbatively NRQCD, HPQCD, ...

NRQCD Factorization

apply NRQCD to
 Bodwin, Braaten and Lepage 1995
 annihilation decays/inclusive production of quarkonium

• motivation:

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infrared divergences for P-waves in CSM decays \chi_{cJ} \rightarrow q \bar{q} g production b \rightarrow \chi_{cJ} + s + g
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- use effective field theory NRQCD to separate hard momentum scales (m_Q and larger) from soft momentum scales ($m_Q v$ and smaller)
- annihilation/creation of QQ pair: hard
- evolution/formation of quarkonium: soft

Annihilation decay rate of charmonium *H*

$$\Gamma[H] = \sum_{n} \hat{\Gamma}[c\bar{c}(n)] \langle H|\mathcal{O}_n|H\rangle$$

- sum over color/angular momentum channels 1 or 8 ¹S₀, ³S₁, ¹P₁, ³P₀, ³P₁, ³P₂, ...
- hard factors: annihilation rate for $c\bar{c}$ into partons expand in powers of $\alpha_s(m_c)$
- soft factors: NRQCD matrix elements scale as powers of *v*
- <u>rigorous</u> factorization formula double expansion in $\alpha_s(m_c)$ and v

Annihilation decay rate of charmonium *H*

$$\Gamma[H] = \sum_{n} \hat{\Gamma}[c\bar{c}(n)] \langle H|\mathcal{O}_n|H\rangle$$

velocity scaling of NRQCD matrix elements

$$\frac{J/\psi:}{\langle \underline{1} \, {}^{3}S_{1} \rangle} \sim v^{3}$$

$$\frac{\langle \underline{8} \, {}^{3}P_{J} \rangle, \langle \underline{8} \, {}^{1}S_{0} \rangle, \langle \underline{8} \, {}^{3}S_{1} \rangle}{\langle \underline{1} \, {}^{3}P_{J} \rangle, \langle \underline{8} \, {}^{3}S_{1} \rangle} \sim v^{7}$$

$$\chi_{cJ}: \qquad (\underline{1} \, {}^{3}P_{J} \rangle, \langle \underline{8} \, {}^{3}S_{1} \rangle \sim v^{5}$$

- solves infrared divergence problem for P-waves
- spin symmetry relates J/ψ , η_c χ_{c0} , χ_{c1} , χ_{c2} , h_c

Inclusive production of charmonium *H*

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

- sum over color/angular momentum channels 1 or 8 ¹S₀, ³S₁, ¹P₁, ³P₀, ³P₁, ³P₂, ...
- hard factors: parton cross sections for creating $c\bar{c}$ expand in powers of $\alpha_s(m_c)$
- soft factors: NRQCD matrix elements scale as powers of *v*
- <u>conjectured</u> factorization formula motivated by <u>perturbative QCD</u> factorization theorems

Inclusive production of charmonium *H*

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

velocity scaling of NRQCD matrix elements

CSM

$$J/\psi: \qquad (\underline{1} \, {}^3S_1) \sim v^3$$

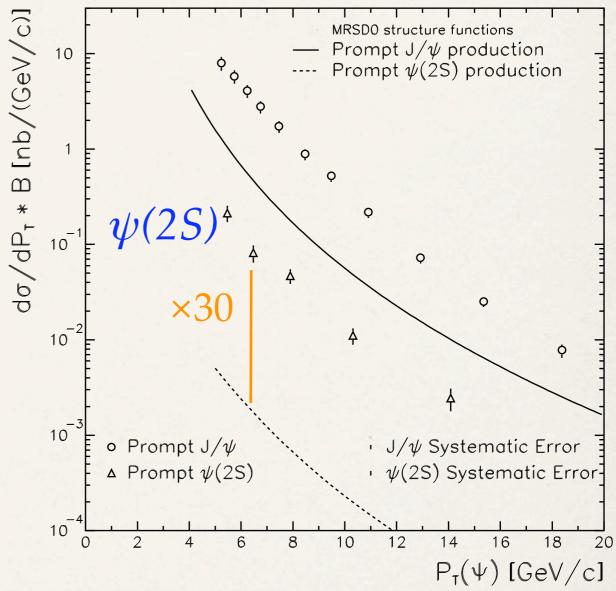
$$\chi_{cJ}: \qquad (\underline{1}^{3}P_{J}), (\underline{8}^{3}S_{1}) \sim v^{5}$$

- solves infrared divergence problem for P-waves
- vacuum saturation approximation
 relates CSM matrix elements for production and decay

$\psi(2S)$ Surplus at the Tevatron

CDF collaboration

• prompt $\psi(2S)$ is 30 times larger than CSM prediction



feeddown from P-wave or D-wave charmonium? charmonium hybrids?

Color evaporation model?

$\psi(2S)$ Surplus at the Tevatron

NRQCD Factorization predicts
 both color-singlet and color-octet production mechanisms

$$\psi(2S): \qquad \langle \underline{1} \, {}^{3}S_{1} \rangle \sim v^{3}$$

$$\langle \underline{8} \, {}^{3}P_{J} \rangle, \langle \underline{8} \, {}^{1}S_{0} \rangle, \langle \underline{8} \, {}^{3}S_{1} \rangle \sim v^{7}$$

• at large p_T ,

CSM term $\langle \underline{1} \, {}^3S_1 \rangle$ is suppressed by α_s^2 color-octet terms $\langle \underline{8} \, {}^1S_0 \rangle$, $\langle \underline{8} \, {}^3P_I \rangle$ are suppressed by $\alpha_s \, v^4$

color-octet term $\langle \underline{8} \, {}^{3}S_{1} \rangle$ is suppressed only by v^{4}

• proposed solution to $\psi(2S)$ surplus: prompt $\psi(2S)$ at large p_T at the Tevatron is dominated by $\langle \underline{8} \, {}^3S_1 \rangle$ term (color-octet mechanism) Braaten and Fleming 1995

Inclusive production of charmonium *H*

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

• for S-waves, truncate after order v^7

$$J/\psi: \qquad \langle \underline{1} \, {}^{3}S_{1} \rangle \sim v^{3}$$

$$\langle \underline{8} \, {}^{3}P_{J} \rangle, \langle \underline{8} \, {}^{1}S_{0} \rangle, \langle \underline{8} \, {}^{3}S_{1} \rangle \sim v^{7}$$

- \Rightarrow 4 universal constants for J/ψ , η_c (1 determined by $J/\psi \rightarrow e^+e^-$)
- for P-waves, truncate after order v^5

$$\chi_{cJ}: \langle \underline{1}^{3}P_{J}\rangle, \langle \underline{8}^{3}S_{1}\rangle \sim v^{5}$$

 \Rightarrow 2 universal constants for χ_{c0} , χ_{c1} , χ_{c2} , h_c (1 determined by $\chi_{c0} \rightarrow \gamma \gamma$)

Inclusive production of charmonium *H*

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

• for S-waves, truncate after order v^7

$$J/\psi: \qquad (\underline{1} \, {}^{3}S_{1}) \sim v^{3}$$

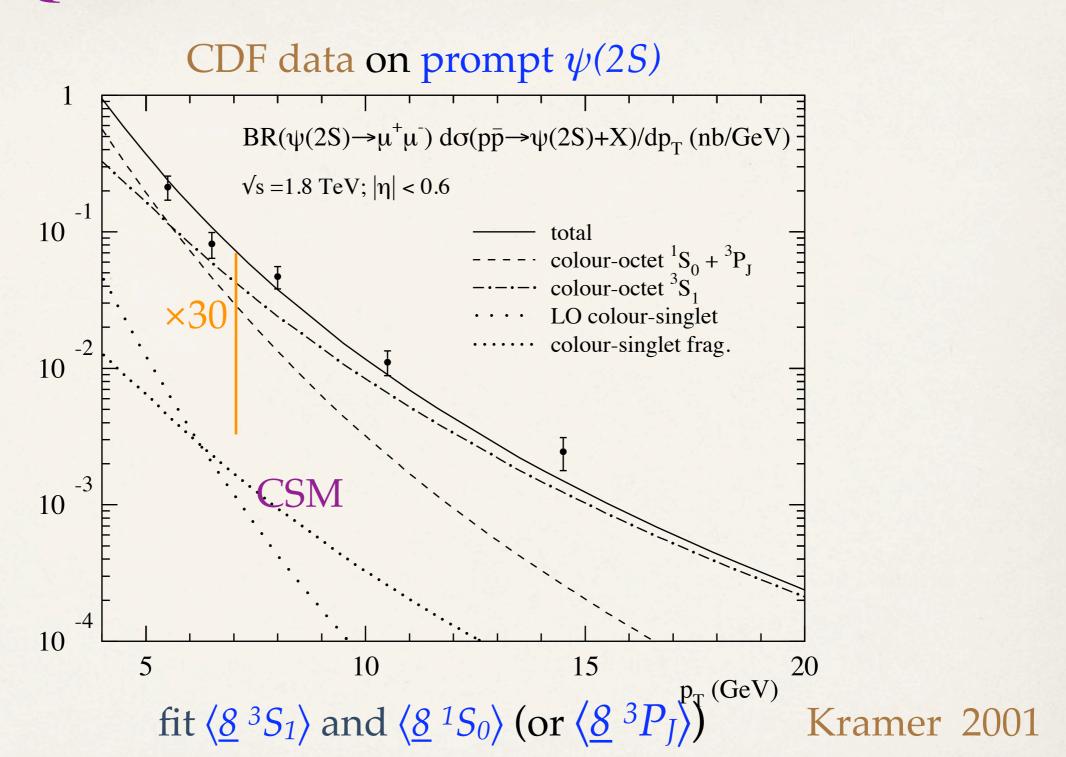
$$\langle \underline{8} \, {}^{3}P_{J}\rangle, (\underline{8} \, {}^{1}S_{0}\rangle, \langle \underline{8} \, {}^{3}S_{1}\rangle \sim v^{7}$$

 \Rightarrow 4 universal constants for J/ψ , η_c (1 determined by $J/\psi \rightarrow e^+e^-$)

• for P-waves, truncate after order v^5

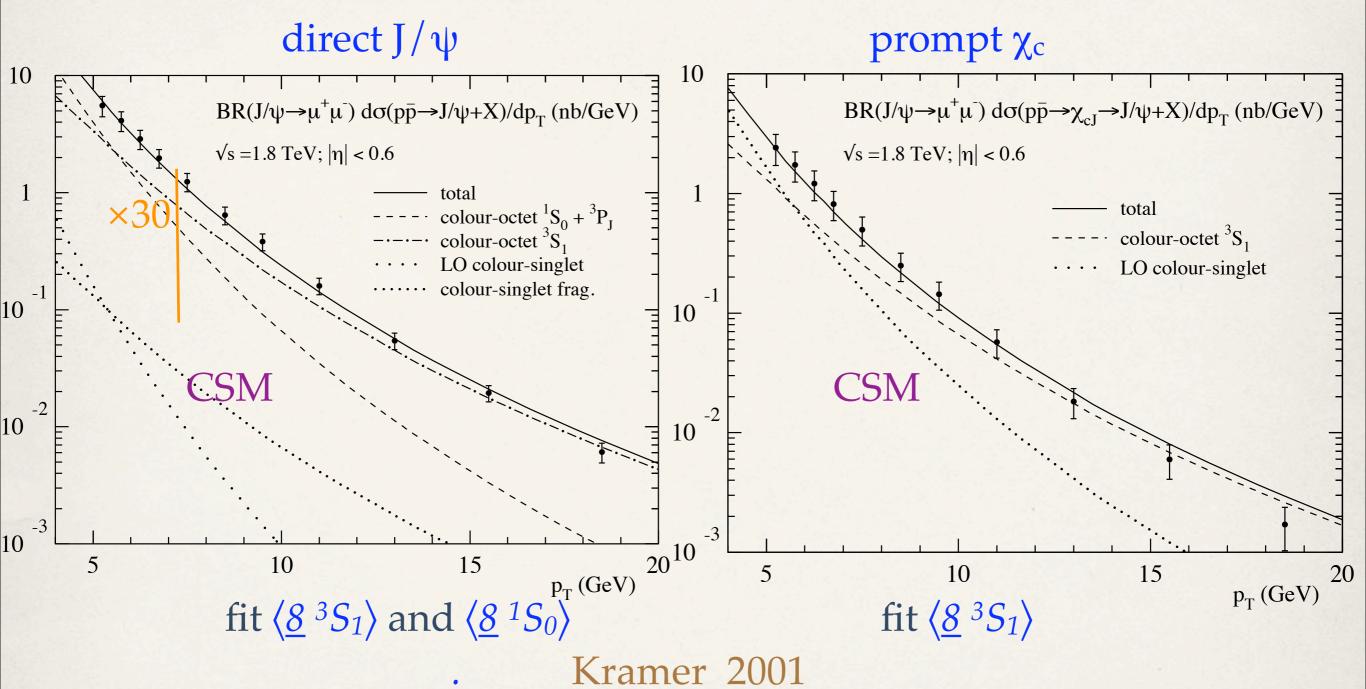
$$\chi_{cJ}: \langle \underline{1}^{3}P_{J} \rangle, \langle \underline{8}^{3}S_{1} \rangle \sim v^{5}$$

 \Rightarrow 2 universal constants for χ_{c0} , χ_{c1} , χ_{c2} , h_c (1 determined by $\chi_{c2} \rightarrow \gamma \gamma$)



NRQCD factorization can accomodate the Tevatron data natural explanation with sound theoretical basis

CDF data



NRQCD factorization can accomodate the Tevatron data natural explanation with sound theoretical basis

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NRQCD Factorization: theory of quarkonium production?

proof of NRQCD factorization formula?

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

- NRQCD matrix elements
 are they universal?
 can they be calculated using lattice QCD?
 is truncation of NRQCD factorization model adequate
- parton cross sections can higher orders in α_s be calculated?

Proof of NRQCD Factorization?

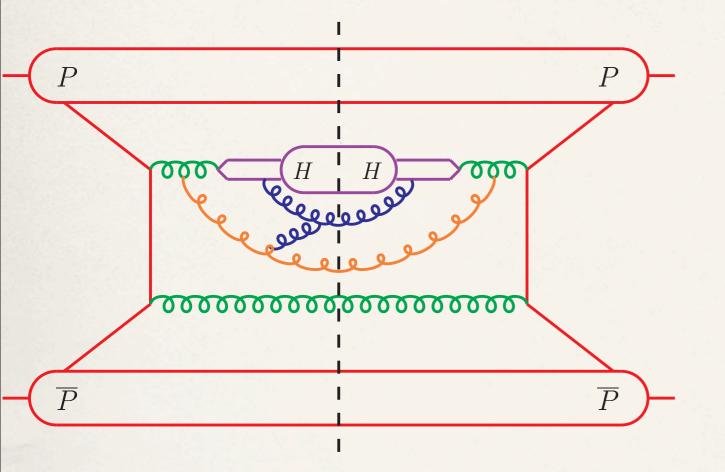
- NRQCD factorization formulas are <u>conjectured</u> motivated by <u>perturbative QCD</u> factorization theorems must be <u>proven</u> to all orders in α_s
- exclusive production

 e⁺ e⁻ → quarkonium + quarkonium
 B → (light meson) + quarkonium
 proof of factorization to all orders in α_s

Bodwin, Lee, Tormo 2008, 2010

Proof of NRQCD Factorization? (cont.) Bodwin @ KITPC

 \bullet Nayak, Qiu, Sterman (2005, 2006): A key difficulty in proving factorization to all orders is the treatment of gluons with momenta of order m_c in the quarkonium rest frame.



- ullet If the orange gluon has momentum of order m_c , it can't be absorbed into the NRQCD matrix element as a quarkonium constituent.
- But the orange gluon can have non-vanishing soft exchanges with the quarkonium constituents.
- The orange gluon can be treated as the eikonal-line part of the NRQCD matrix element, provided that the answer does not depend on the direction of the eikonal line (universality of the matrix elements).
- Nayak, Qiu, Sterman (2005, 2006): At two-loop order, the eikonal lines contribute but a "miracle" occurs: The dependence on the direction of the eikonal line cancels.
- In general, factorization of the inclusive cross section beyond two-loop order is still an open question.

 \bullet An all-orders proof is essential because the α_s associated with soft gluons is not small.

Proof of NRQCD Factorization? (cont.)

dramatic new development at this workshop!

proof of factorization to all orders in α_s at large p_T !

Jian-Wei Qiu and collaborators

- separate very hard scale p_T from not-so-hard scale m_Q by expanding in powers of m_Q^2/p_T^2
- at leading power: factorization (parton fragmentation) at order m_Q^2/p_T^2 : factorization (QQ fragmentation) at higher orders: no factorization?
- not-so-hard factor involves scale m_Q and softer scales $m_Q v$ and smaller use NRQCD factorization to express it in terms of NRQCD matrix elements?

NRQCD matrix elements

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

- are they universal?
 in absence of proof, use phenomenology
- can they be calculated using lattice QCD? CSM matrix elements: YES, up to $O(v^4)$

color-octet matrix elements: NO

truncation of NRQCD factorization model

S-waves: $\langle \underline{1} \, {}^{3}S_{1} \rangle$, $\langle \underline{8} \, {}^{3}S_{1} \rangle$, $\langle \underline{8} \, {}^{1}S_{0} \rangle$, $\langle \underline{8} \, {}^{3}P_{I} \rangle$

P-waves: $\langle \underline{1} \, {}^{3}P_{J} \rangle$, $\langle \underline{8} \, {}^{3}S_{1} \rangle$

is this sufficiently accurate for charmonium? maybe for bottomonium? maybe

Parton cross sections

$$d\sigma[H] = \sum_{n} d\hat{\sigma}[c\bar{c}(n)] \langle \mathcal{O}_{n}^{H} \rangle$$

accurate predictions require at least NLO in α_s

for charmonium, $\alpha_s(m_c) \approx 0.25$

for bottomonium, $\alpha_s(m_b) \approx 0.18$

photoproduction
 Artoisenet, Campbell, Maltoni, Tramontano 2009
 Chang, Li, Wang 2009; Li, Chao 2009
 Butenschoen, Kniehl 2009

• yy collisions Klasen, Kniehl, Mihaila, Steinhauser 2005

Phang, Gao, Chao 2005; Zhang, Ma, Chao 2008

• e⁺ e⁻ → double charmonium Gong, Wang 2008

Zhang, Chao 2006; Ma, Zhang, Chao 2008 Gong, Wang 2008, 2009

• e⁺ e⁻ → charmonium + X Zhang, Ma, Wang, Chao 2009

Petrelli, Cacciari, Greco, Maltoni, Mangano 1988

Campbell, Maltoni, Tramontano 2008; Artoisenet, Lansberg, Maltoni, 2008

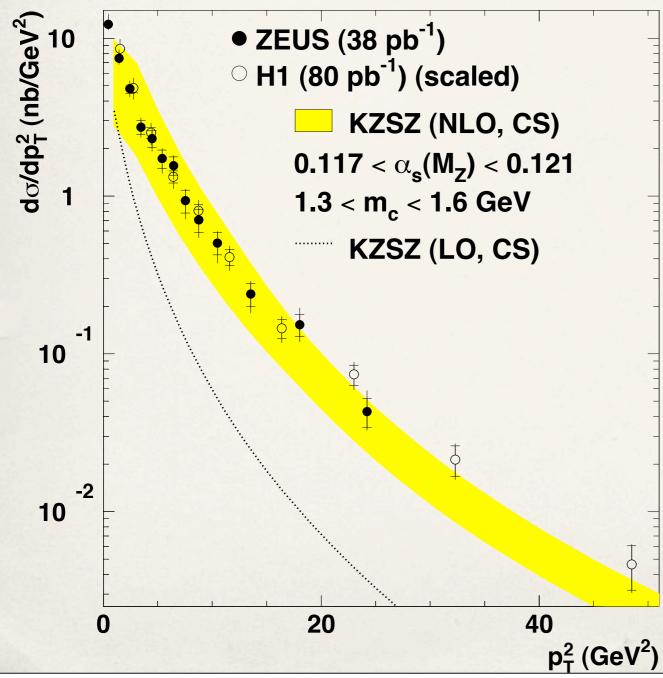
• hadron collisions Li, Wang 2008; Gong, Wang 2008; Gong, Li, Wang 2009

Phenomenological Status of NRQCD Factorization

- photoproduction
- $e^+e^- \rightarrow$ double charmonium
- $e^+e^- \rightarrow charmonium + X$
- γγ collisions
- hadron collisions

Inelastic J/ψ Photoproduction Cross Section at HERA

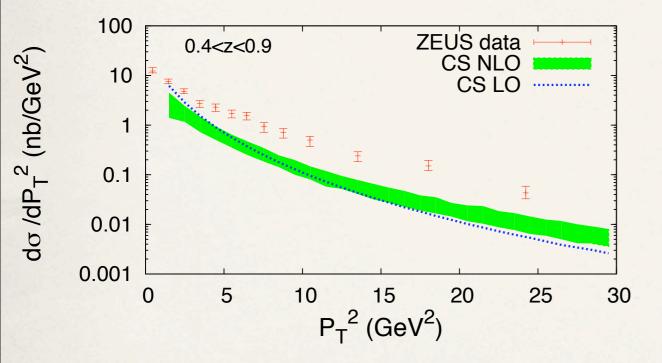
 It had been believed that NLO color-singlet calculations leave little room for a color-octet contribution.

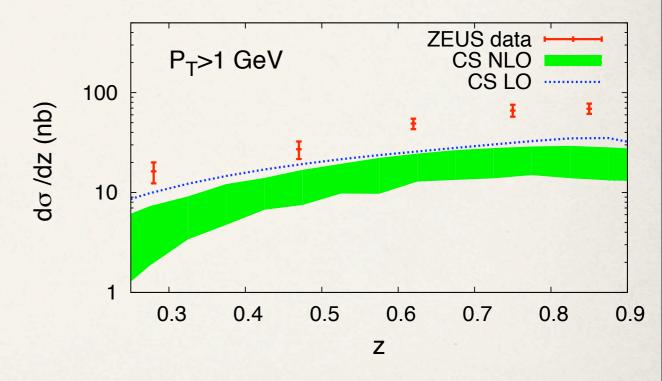


- NLO corrections increase the colorsinglet contribution substantially. (Krämer, Zunft, Steegborn, Zerwas (1994); Krämer (1995))
- NLO corrections include $\gamma + g \rightarrow (c\bar{c}) + gg$, which is dominated by t-channel gluon exchange.
- For large p_T , this process goes as $\alpha_s^3 m_c^2/p_T^6$, instead of $\alpha_s^2 m_c^4/p_T^8$.

Recent Theoretical Developments

- Artoisenet, Campbell, Maltoni, Tramontano (2009): A new calculation of NLO color-singlet contribution
 - Confirms the analytic results of previous calculations.
 - But a more reasonable choice of renormalization/factorization scale $(\sqrt{4m_c^2+p_T^2} \text{ instead of } m_c/\sqrt{2})$ yields much smaller numerical results for cross sections.





- Leaves room for a color-octet contribution.
- There is no longer an obvious conflict between the NRQCD prediction and the HERA data.

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Exclusive Double-Charmonium Production at Belle and BABAR

$$e^+e^- \rightarrow J/\psi + \eta_c$$

Experiment

Belle (2004):
$$\sigma[e^+e^- \to J/\psi + \eta_c] \times B_{>2} = 25.6 \pm 2.8 \pm 3.4 \text{ fb.}$$

BABAR (2005): $\sigma[e^+e^- \to J/\psi + \eta_c] \times B_{>2} = 17.6 \pm 2.8^{+1.5}_{-2.1} \text{ fb.}$

larger by ≈5

ullet NRQCD at LO in $lpha_s$ and v

Braaten, Lee (2003):
$$\sigma[e^+e^- \to J/\psi + \eta_c] = 3.78 \pm 1.26$$
 fb.

Liu, He, Chao (2003):
$$\sigma[e^+e^- \rightarrow J/\psi + \eta_c] = 5.5$$
 fb.

The two calculations employ different choices of m_c , NRQCD matrix elements, and α_s .

Braaten and Lee include QED effects.

Confirmed by Brodsky, Ji, and Lee in light-front QCD in the quarkonium nonrelativistic limit.

- Exclusive process: the color-octet contribution is suppressed as v^4 .
- The LO color-singlet matrix elements are determined from $\eta_c \to \gamma \gamma$ and $J/\psi \to e^+e^-$.

$$\alpha_s$$
 Corrections to $e^+e^- \rightarrow J/\psi + \eta_c$

- An important step in resolving the discrepancy:
 - Zhang, Gao, Chao (2005) found that corrections at NLO in α_s yield a K factor of about 1.96.
- Confirmed by Gong and Wang (2007).
- Not enough to bring theory into agreement with experiment.

relativistic corrections: ×1.5?

- Theory and experiment agree within uncertainties:
 - Theory: $\sigma[e^+e^- \to J/\psi + \eta_c] = 17.6^{+8.1}_{-6.7} \, \mathrm{fb}$
 - Belle: $\sigma[e^+e^- \to J/\psi + \eta_c] \times B_{>2} = 25.6 \pm 2.8 \pm 3.4 \text{ fb.}$
 - BABAR: $\sigma[e^+e^- \to J/\psi + \eta_c] \times B_{>2} = 17.6 \pm 2.8^{+1.5}_{-2.1}$ fb.
- Caveat: $B_{>2}$ is not known.
 - Could be as small as 0.5-0.6.
 - Even so, the error bars of theory and the BABAR experiment overlap.
- Zhang, Ma, Chao (2008): In the cases of $\sigma[e^+e^- \to J/\psi(\psi(2S)) + \chi_{c0}]$, large K factors (~ 2.8) may bring theory into agreement with experiment.

Inclusive Double $c\bar{c}$ Production at Belle

• Belle (2002):

$$\frac{\sigma[e^+e^- \to J/\psi + c\bar{c} + X]}{\sigma[e^+e^- \to J/\psi + X]} = 0.59^{+0.15}_{-0.13} \pm 0.12$$

pQCD plus color-singlet model (Cho, Leibovich (1996); Baek, Ko, Lee, Song (1997); Yuan, Qiao,
 Chao (1997)):

$$\frac{\sigma[e^+e^- \to J/\psi + c\bar{c} + X]}{\sigma[e^+e^- \to J/\psi + X]} \approx 0.1$$

• There is a significant disagreement between experiment and the LO color-singlet model.

Inclusive I/ψ + charm from e^+e^- (cont.) Bodwin @ KITPC

NLO corrections: Zhang, Chao 2007; Gong, Wang 2009 Ma, Zhang, and Chao 2008; Gong, Wang 2009

Effect of NLO calculations on the ratio

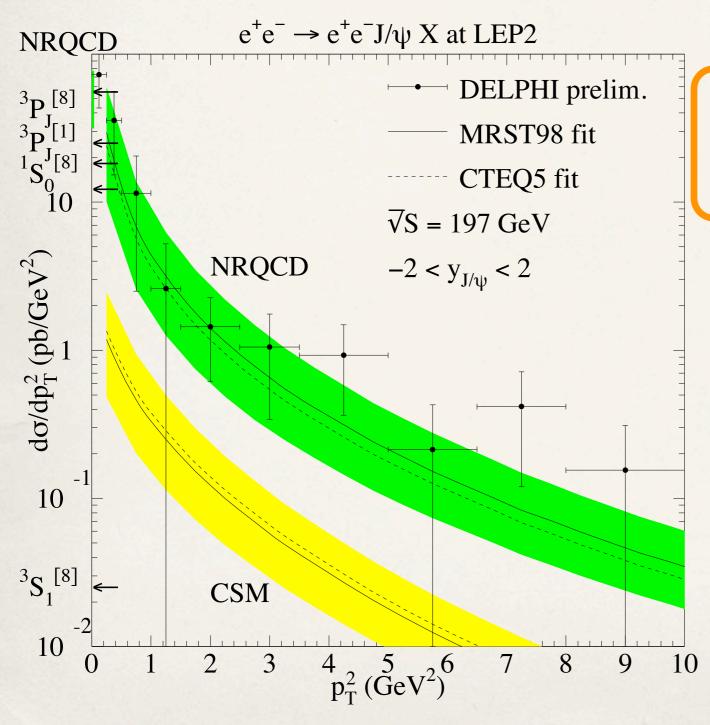
 NLO calculations significantly reduce the discrepancy between theory and experiment for the ratio of cross sections:

$$\frac{\sigma[e^+e^- \to J/\psi + c\bar{c} + X]}{\sigma[e^+e^- \to J/\psi + X]} \approx 0.5$$

- Only color-singlet contributions are included.
- No longer an apparent disagreement between experiment and color-singlet theory.
- It would be good to have a detailed error analysis for the theoretical prediction.
- ullet It is important for BABAR to check the Belle results for inclusive double- $car{c}$ production.

Belle:
$$\frac{\sigma[e^+e^- \to J/\psi + c\bar{c} + X]}{\sigma[e^+e^- \to J/\psi + X]} = 0.59^{+0.15}_{-0.13} \pm 0.12$$

$$\gamma\gamma \to J/\psi + X$$
 at LEP



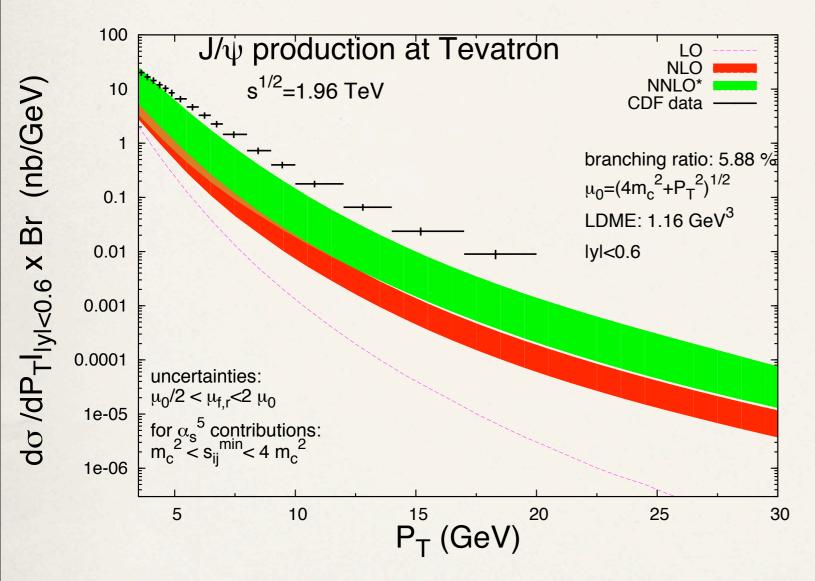
- Comparison of theory (Klasen, Kniehl, Mihaila, Steinhauser) with Delphi data clearly favors NRQCD over the color-singlet model.
- Theory uses Braaten-Kniehl-Lee matrix elements from Tevatron data and MRST98LO (solid) and CTEQ5L (dashed) PDF's.
- Theoretical uncertainties from
 - Renormalization and factorization scales (varied by a factor 2),
 - NRQCD color-octet matrix elements,
 - Different linear combination of matrix elements than in Tevatron cross sections.

J/ψ Production in DIS at HERA

- Note that NLO calculations are not yet available for this process.
- The NRQCD (Kniehl, Zwirner (2001)) prediction uses Braaten-Kniehl-Lee (1999) matrix elements extracted from the Tevatron data and MRST98LO and CTEQ5L PDF's.
- Theoretical uncertainties from
 - PDF's,
 - Renormalization and factorization scales (varied by a factor 2),
 - NRQCD color-octet matrix elements,
 - Different linear combination of matrix elements than in Tevatron cross sections.

New Results for J/ψ Production

Color-singlet contribution:



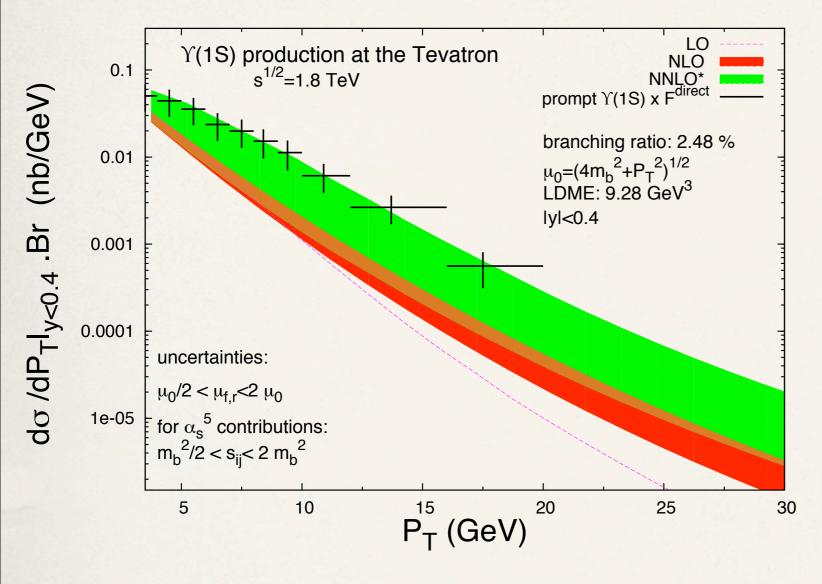
- Plot from Pierre Artoisenet, based on work by Artoisenet, Campbell, Lansberg, Maltoni, Tramontano (in progress)
- The NNLO* calculation is an estimate based on real-emission contributions only.
- The data still seem to require a coloroctet contribution, but its size may be reduced from previous estimates.
 Affects the matrix elements used to

Affects the matrix elements used to compute all other processes.

Color-octet contribution:

NLO corrections are about 14% (Gong, Li, and Wang (2008)).

New Results for Color-Singlet \(\gamma \) Production



- Plot from Pierre Artoisenet, based on work by Artoisenet, Campbell, Lansberg, Maltoni, Tramontano (2008)
- NLO results confirmed by Gong and Wang (2007).

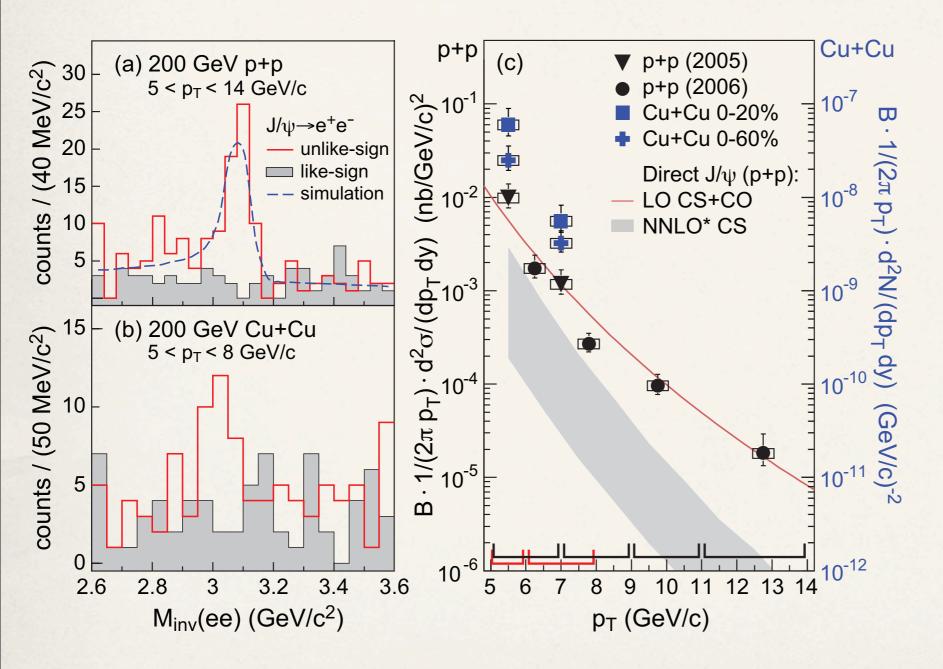
- The data could be explained by color-singlet production alone.
- There is still room for a substantial amount of color-octet production.
- Color-octet production is suppressed as v^4 . Should be smaller for Υ ($v^2 \approx 0.1$) than for J/ψ ($v^2 \approx 0.3$).

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J/ψ Production at RHIC

• The STAR Collaboration has measured the J/ψ p_T distributions in p+p and Cu+Cu collisions:



- The LO color-singlet plus color-octet calculation (Nayak, Liu, Cooper (2003)) fits the data well.
 - Does not include feeddown from $\psi(2S)$, χ_c , or B decays. (Estimated to be a factor 1.5.)

Polarization

NRQCD factorization

predicts the polarization of quarkonium with no additional parameters

dramatic qualitative prediction for hadron collisions: direct J/ψ , Y transversely polarised at large p_T Cho, Wise 1995

• at sufficiently large p_T , charmonium production is dominated by gluon fragmentation

$$g + g \rightarrow g^* + g$$

• at LO in α_s , gluon fragments into color-octet $c\overline{c}$ pair that inherits transverse polarization of gluon

$$g^* \uparrow \rightarrow \bar{cc}(\underline{8} \, {}^3S_1) \uparrow$$

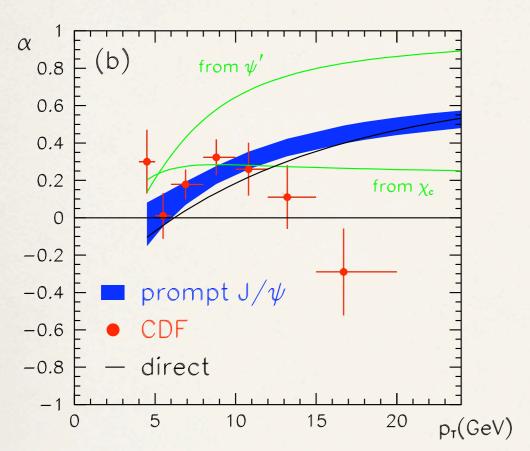
• at LO in v, hadronization into 3S_1 charmonium preserves transverse polarization of $c\bar{c}$ pair

$$c\bar{c}(\underline{8}^{3}S_{1})\uparrow \rightarrow J/\psi\uparrow + X$$

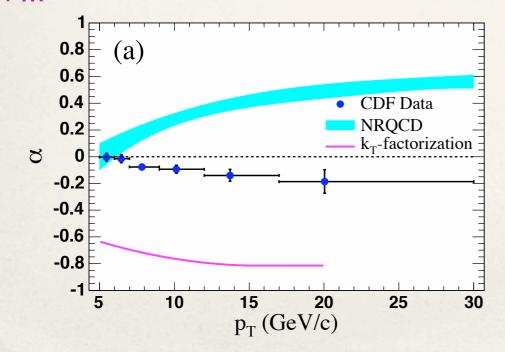
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J/ψ Polarization

Run I:

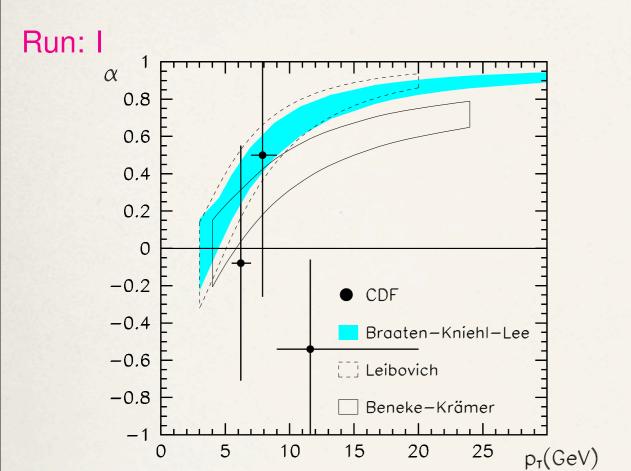


Run II:

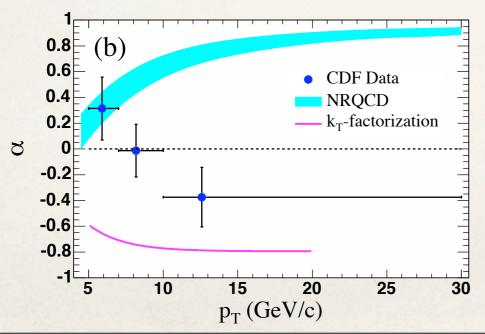


- $d\sigma/d(\cos\theta) \propto 1 + \alpha \cos^2\theta$.
 - $-\alpha = 1$ is completely transverse;
 - $-\alpha = -1$ is completely longitudinal.
- NRQCD prediction from Braaten, Kniehl, Lee (1999).
 - Feeddown from χ_c states is about 30% of the J/ψ sample and dilutes the polarization.
 - Feeddown from $\psi(2S)$ is about 10% of the J/ψ sample and is largely transversely polarized.
- Run I results are marginally compatible with the NRQCD prediction.
- Run II results are inconsistent with the NRQCD prediction.
- Also, inconsistent with Run I results.
 CDF was unable to track down the source of the Run I-Run II discrepancy.

$\psi(2S)$ Polarization

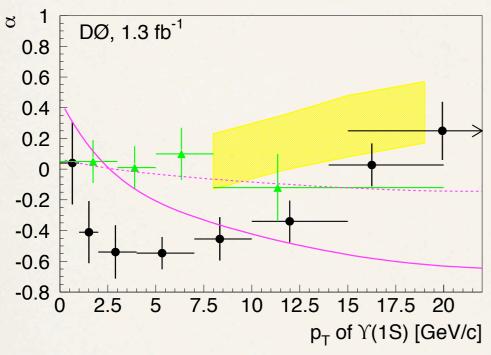


Run: II

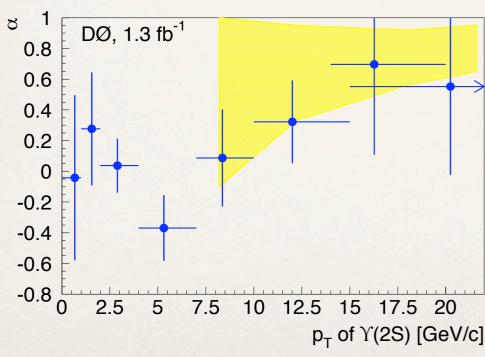


 The Run II data are incompatible with the NRQCD prediction.

$\Upsilon(1S)$ Polarization:



$\Upsilon(2S)$ Polarization:



- In the $\Upsilon(1S)$ case, the D0 results (black) are incompatible with the CDF results (green).
- The CDF results are compatible with the NRQCD prediction (yellow).
- The D0 results are marginally incompatible with the NRQCD prediction.
- ullet The curves are the limiting cases of the k_T -factorization prediction.
- In the $\Upsilon(2S)$ case, the theoretical and experimental error bars are too large to make a stringent test.

Polarization (cont.)

Are still higher order perturbative QCD calculations necessary for polarization?

Does NRQCD factorization fail for polarization?

new development at this workshop! large- p_T factorization Jian-Wei Qiu et al.

- expand in powers of m_Q^2/p_T^2 to separate scales m_Q and p_T
- at leading power: parton fragmentation \Rightarrow T at order m_Q^2/p_T^2 : QQ fragmentation \Rightarrow L
- Will predictions including QQ fragmentation agree with data?

Stumbling towards a Theory of Quarkonium Production

Color-singlet model (1976-1995) Color evaporation model (1977-?)



NRQCD factorization

- still a viable theory of quarkonium production!
- exclusive quarkonium: proven to all orders
- inclusive quarkonium: verified to NNLO
- can it be combined with large- p_T factorization?

NRQCD factorization model

S-wave multiplets: 3 color-octet parameters

P-wave multiplets: 1 color-octet parameter

• still a viable model of charmonium production bottomonium production

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Stumbling towards a Theory of Quarkonium Production

NLO perturbative QCD corrections

- have removed most dramatic discrepancies between NRQCD factorization and experiment (polarization is important exception)
- decrease the importance of color-octet contributions

Large- p_T factorization

Jian-Wei Qiu et al.

- separates scales m_Q and p_T
- introduces QQ fragmentation
- still requires NRQCD factorization to reduce production rates to a few constants

Stumbling towards a Theory of Quarkonium Production Experimental outlook

- final results from B factories (Belle, Babar)

 DESY (H1, Zeus)

 Tevatron (CDF, D0)
- first results from LHC experiments
 extend charmonium out to fragmentation region
 high statistics measurements of bottomonium
 additional results from RHIC
- future results from super-B factories

Will NRQCD factorization

remain a viable theory of quarkonium production?